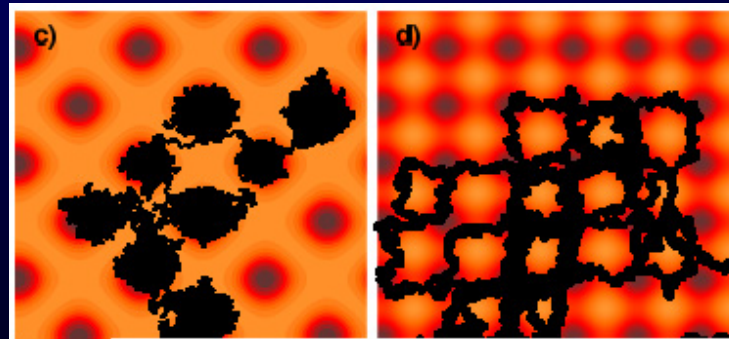
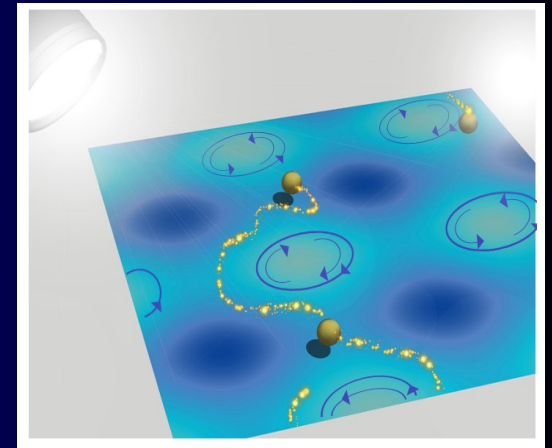


# Giant Enhanced Diffusion of Nanoparticles in Optical Vortex Fields

*Juan José Sáenz*

MoLE ("Moving Light & Electrons") group  
([www.uam.es/mole](http://www.uam.es/mole))

[Universidad Autónoma de Madrid]



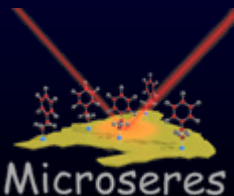
*Silvia Albaladejo*  
*Manuel I. Marqués*  
*Juanjo Sáenz* (UAM-Mole, Madrid, Spain)

*Marine Laroche* (Institute d'Optique, CNRS,  
Paris, France)  
*Frank Scheffold* (Fribourg Univ., Switzerland)

*Ivar Zapata*  
*Jose Maria R. Parrondo*  
*Fernando Sols* (Univ. Complutense Madrid,  
Spain)

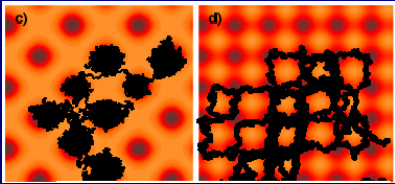
(\*) *Donostia International Physics Center (DIPC)*

Phys. Rev. Lett. **102**, 113602 (2009)  
Phys. Rev. Lett. **103**, 130601 (2009)  
NanoLett. **9**, 3527 (2009)

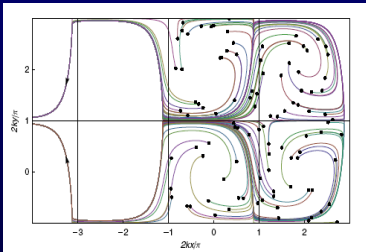




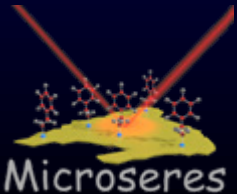
- Introduction
  - Optical Forces on small particles.



- Giant enhanced diffusion of gold nanoparticles on optical vortex lattices

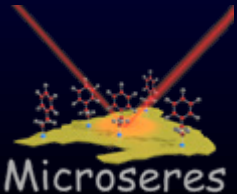


- Deterministic Ratchet from Stationary Light vortex Fields.

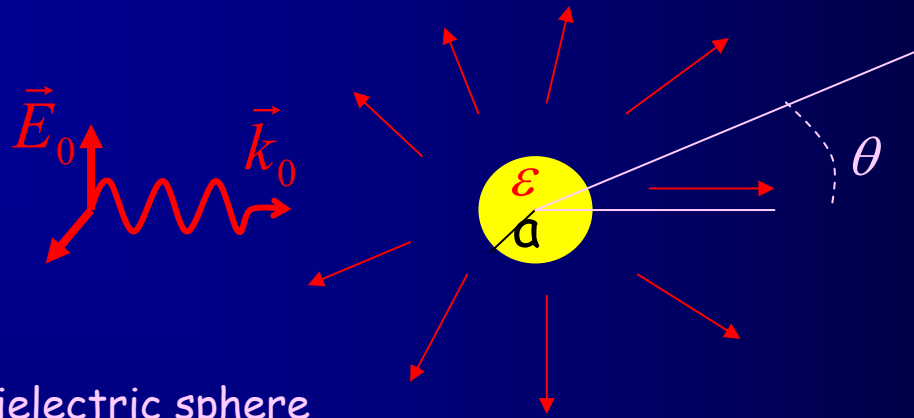


# Scattering from a small particle

- Electric dipole
- Polarizability



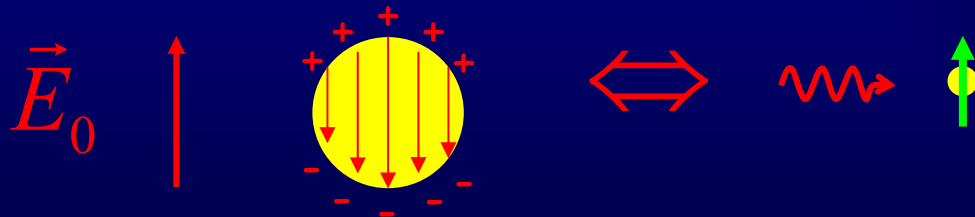
# Scattering by a small particle



Small dielectric sphere  
(Classical Electrodynamics, Jackson)

$$\lambda \gg a$$

$$\vec{p} = \epsilon_0 \epsilon_h \alpha \vec{E}_0$$



Within the electrostatic approximation,  $\alpha$  is real.

$$\alpha_0 = 4\pi a^3 \left( \frac{\epsilon - \epsilon_h}{\epsilon + 2\epsilon_h} \right)$$

$$\vec{p} = \epsilon_0 \epsilon_h \alpha \vec{E}_0$$

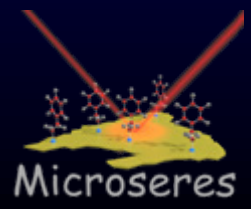
$$\alpha = \frac{\alpha_0}{1 - i \frac{k^3}{6\pi} \alpha_0}$$

$$k \operatorname{Im}\{\alpha\} = \sigma$$

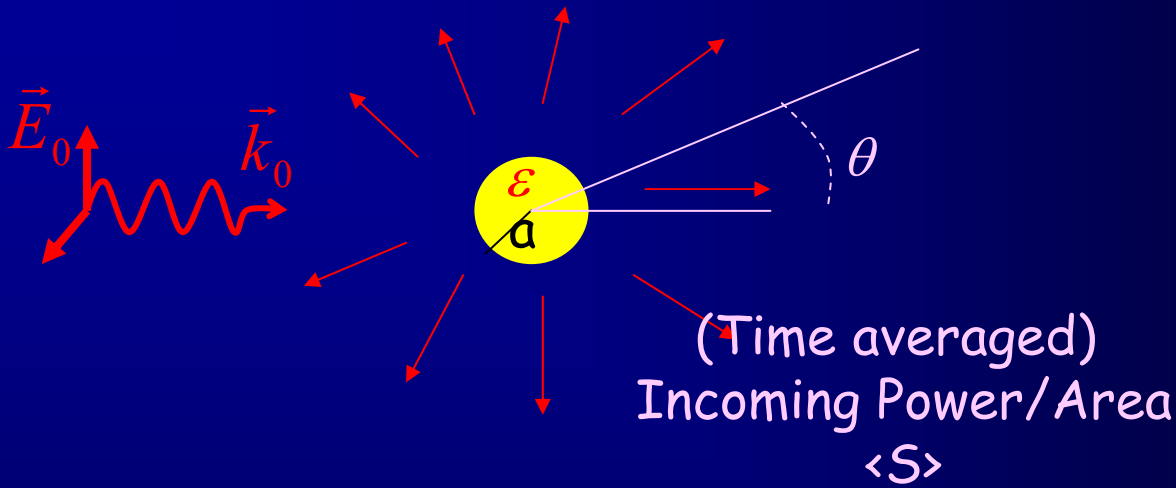
$$\operatorname{Re}\{\alpha\} \approx \alpha_0$$

$$\operatorname{Im}\left\{\frac{1}{\alpha}\right\} = -\frac{k^3}{6\pi} = -k^2 \operatorname{Im}\{G_{ii}(0,0)\}$$

$\lambda \gg a$



Let us consider plane wave incidence

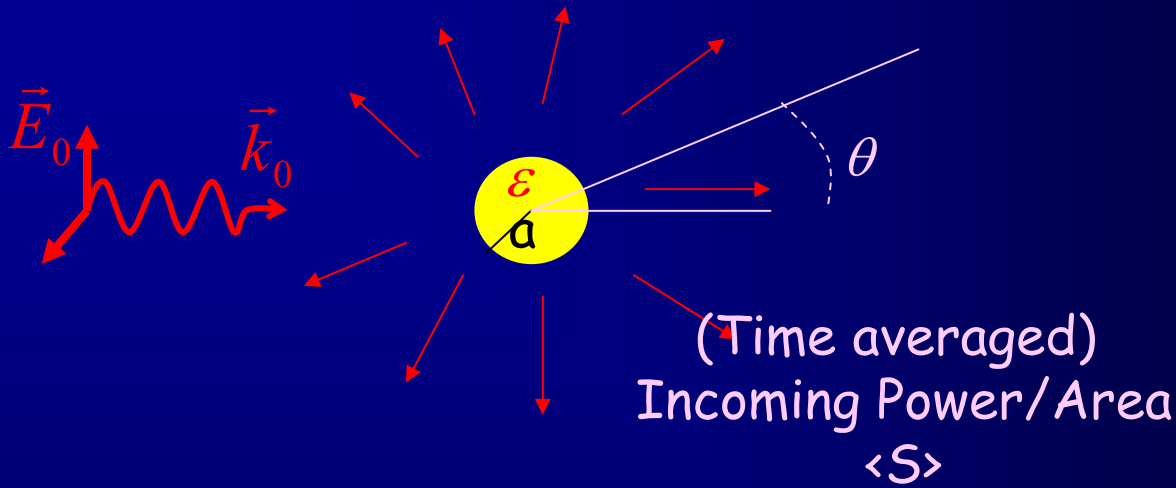


$$\frac{1}{2} \Re \{ \mathbf{E} \times \mathbf{H}^* \}$$

Scattering Cross Section =  
(Scattered Power)/(Inc. Power/area)

$$\Leftrightarrow \sigma_{scat} = k^3 |\alpha|^2 \text{Im} \{ G_{ii}(0,0) \} = \frac{k^4}{6\pi} |\alpha|^2$$

Let us consider plane wave incidence



$$\frac{1}{2} \Re \{ \mathbf{E} \times \mathbf{H}^* \}$$

Extinction Cross Section =  
(Power removed from the beam)/(Inc. Power/area)

$$\sigma_{ext} = k \operatorname{Im} \{ \alpha \}$$



$$\sigma = k \operatorname{Im}\{\alpha\} \quad \text{Extinction Cross Section}$$

In absence of absorption:  
Transmitted + Reflected Power = Incoming Power

$$\operatorname{Im}\{\alpha\} = k^2 |\alpha|^2 \operatorname{Im}\{G_{ii}\}$$

$$\operatorname{Im}\left\{\frac{1}{\alpha}\right\} = -k^2 \operatorname{Im}\{G_{ii}(0,0)\} \Leftrightarrow \sigma = k^3 |\alpha|^2 \operatorname{Im}\{G_{ii}(0,0)\} = \frac{k^4}{6\pi} |\alpha|^2$$

If  $\alpha$  is real:  
Power is not conserved



# Optical Forces

For small particles

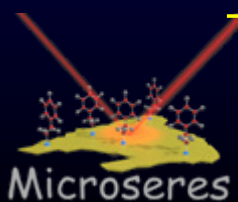
$$\vec{F} = \frac{d}{dt} \vec{P}_{mec} = \int_V (\rho \vec{E} + \vec{J} \times \vec{B}) dV$$

$$\langle \vec{F} \rangle_t = \frac{1}{2} \text{Re} \left\{ \sum_i p_i \vec{\nabla} E_i^* \right\} =$$

$$\langle F \rangle = \text{Re} \{ \alpha \} \left\{ \nabla \frac{1}{4} |E|^2 \right\} + \sigma \frac{1}{2} \text{Re} \left\{ \frac{1}{c} E \times H^* \right\}$$

$$+ \sigma \frac{1}{2} \text{Re} \left\{ i \frac{\epsilon_0}{k_0} (E \cdot \nabla) E^* \right\}$$

????



# Optical Forces (Maxwell eqs. + Lorentz Force)

$$\langle \vec{F} \rangle_t = \frac{1}{2} \text{Re} \left\{ \sum_i p_i \vec{\nabla} E_i^* \right\} =$$

(A. Ashkin, PRL (70))

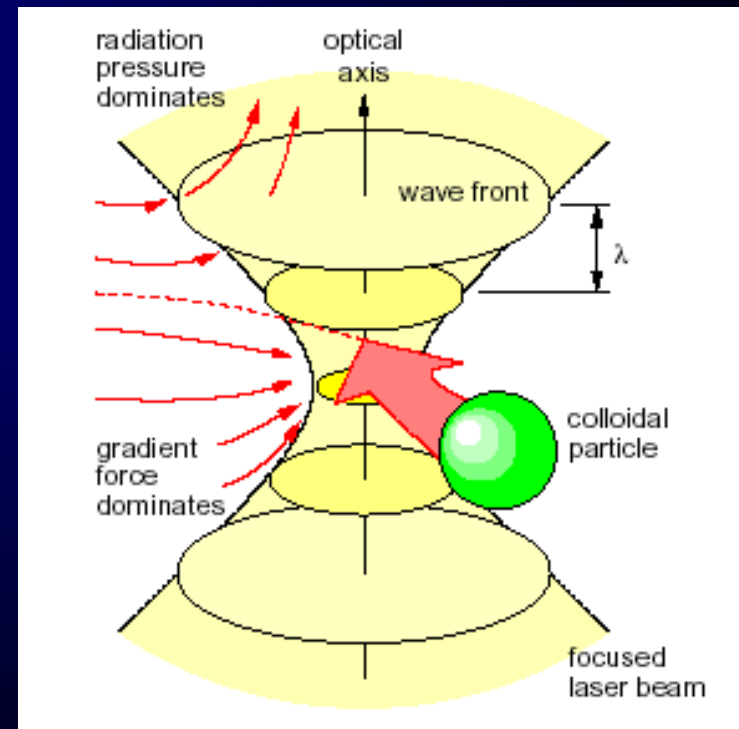
## Optical Tweezers

$$\frac{1}{4} \alpha_0 \vec{\nabla} |\vec{E}_0|^2 \quad \text{Polarization}$$

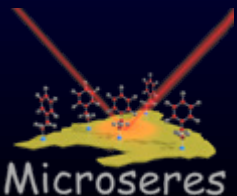
+

$$\frac{1}{c} \langle \vec{S} \rangle \sigma \quad \text{Radiation Pressure}$$

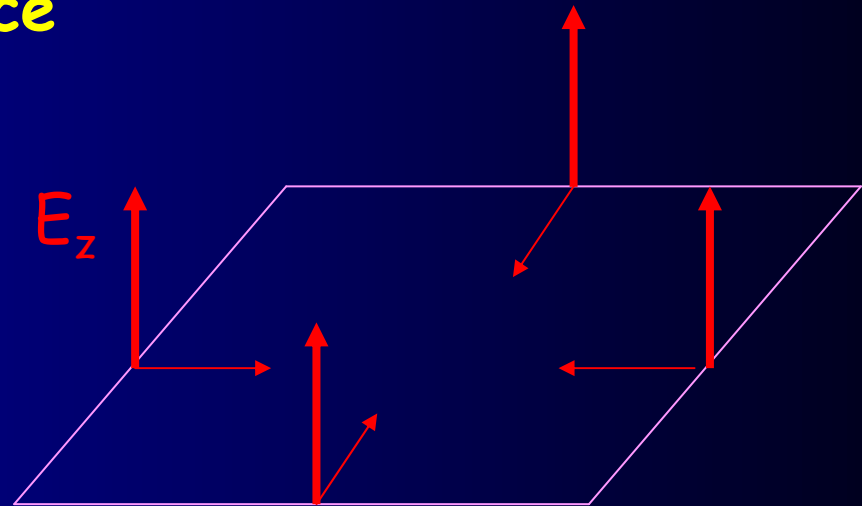
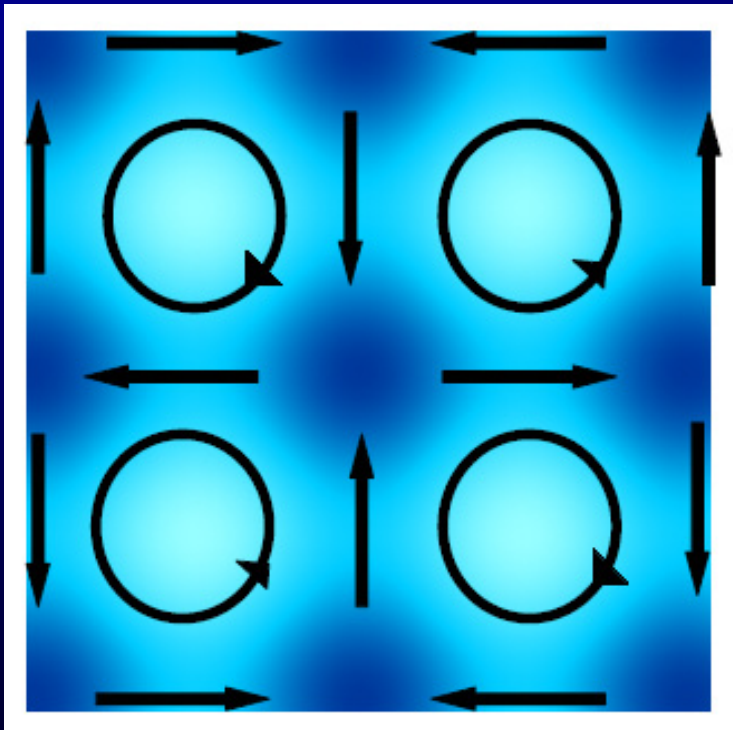
(J.P. Gordon, PRA (73))



D.G. Grier, Nature (03)



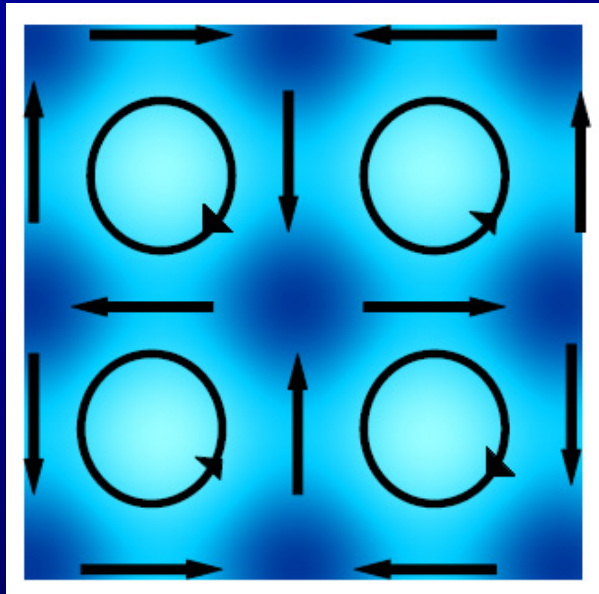
# Optical vortex lattice



A. Hemmerich and T.W. Hänsch,  
Phys. Rev. Lett. 68, 1492 (1992)

$$E_z(x, y; \omega) = \frac{2iE_0}{\sqrt{2}} (\sin k_0x + e^{i\phi} \sin k_0y)$$

# Optical vortex lattice



90°

$$\mathbf{F}^{(s)} = \Re\{\alpha\} \frac{1}{2} \nabla |E_z|^2 + \sigma \left\{ \frac{1}{c} \langle \mathbf{S} \rangle^{(s)} \right\}$$

$$\frac{1}{c} \langle \mathbf{S} \rangle^{(s)} = c \nabla \times \langle L_O \rangle$$

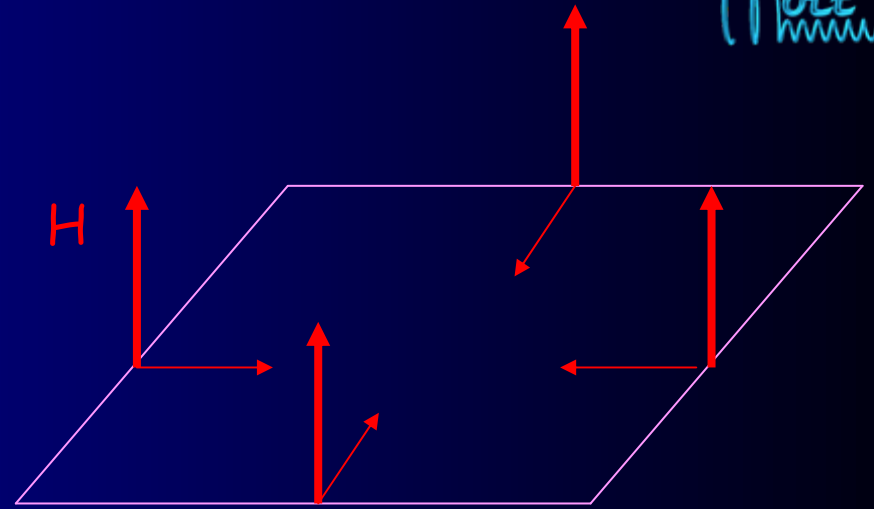
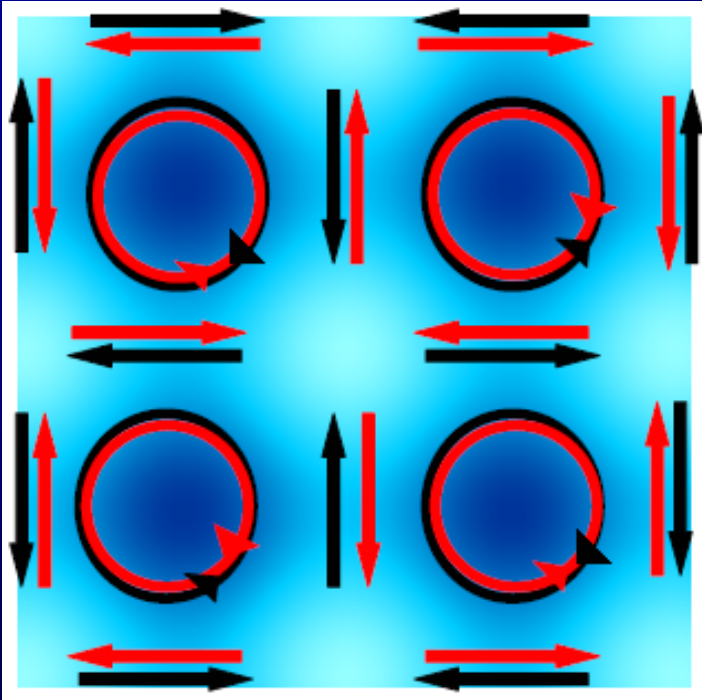
$$\langle L_O \rangle \equiv \frac{\epsilon_0}{\omega} |E_0|^2 \sin \phi \cos k_0 x \cos k_0 y \mathbf{u}_z$$

$L_O \Rightarrow$  Orbital Angular Momentum

$$\frac{|E_z|^2}{|E_0|^2} = \sin^2 k_0 x + \sin^2 k_0 y + 2 \cos \phi \sin k_0 x \sin k_0 y$$

The curl force can be associated to the orbital angular momentum,  $L_O$  arising as a consequence of the rotation of the Poynting vector around the field nodes.

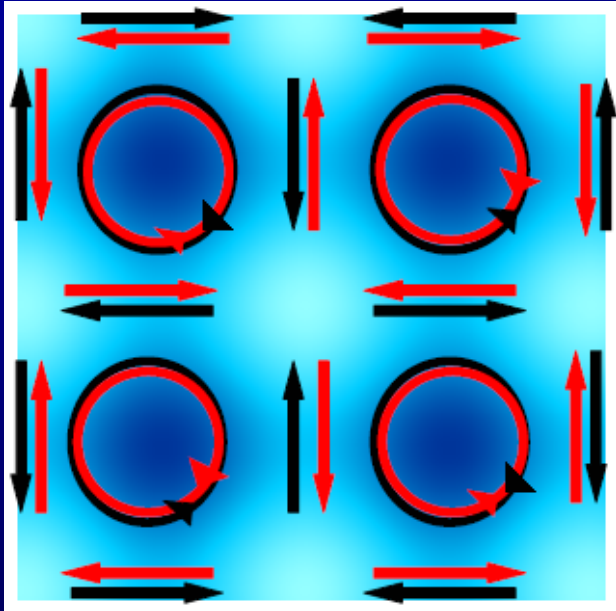
# Optical lattice



A. Hemmerich and T.W. Hänsch,  
Phys. Rev. Lett. 70, 410 (1993)

$$H_z(x, y; \omega) = -\epsilon_0 c \frac{2iE_0}{\sqrt{2}} (\sin k_0 x + e^{i\phi} \sin k_0 y)$$

# Optical lattice



$$H_z(x, y; \omega) = -\epsilon_0 c \frac{2iE_0}{\sqrt{2}} (\sin k_0 x + e^{i\phi} \sin k_0 y)$$

$$E_x = \frac{2E_0}{\sqrt{2}} e^{i\phi} \cos k_0 y$$

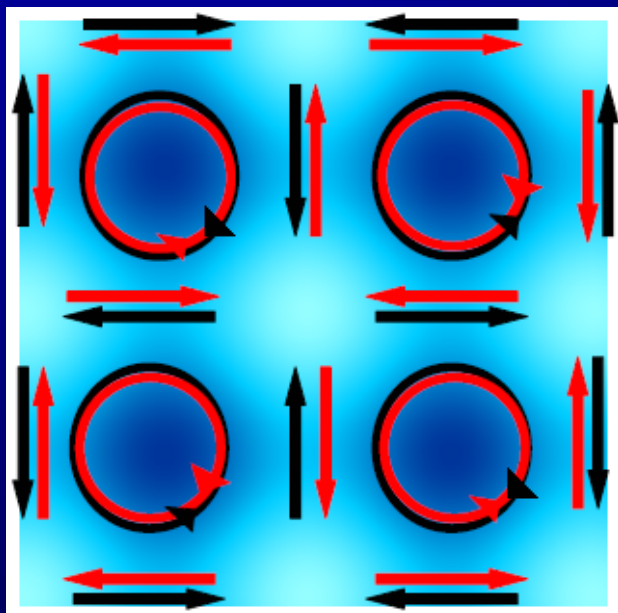
$$E_y = -\frac{2E_0}{\sqrt{2}} \cos k_0 x.$$

$$F^{(p)} = \Re\{\alpha\} \frac{1}{2} |E_0|^2 \nabla (\cos^2 k_0 x + \cos^2 k_0 y).$$

Does not depend on the Phase Shift !

CONSERVATIVE !

## Optical lattice



$$F^{(p)} = \Re\{\alpha\} \frac{1}{2} |E_0|^2 \nabla (\cos^2 k_0 x + \cos^2 k_0 y).$$

Poynting vector?

$$\frac{1}{c} \langle S \rangle^{(s)} = c \nabla \times \langle L_O \rangle$$

$$\langle L_O \rangle \equiv \frac{\epsilon_0}{\omega} |E_0|^2 \sin \phi \cos k_0 x \cos k_0 y \mathbf{u}_z$$



# Optical Forces

(Maxwell eqs. + Lorentz Force)



$$\langle \vec{F} \rangle_t = \frac{1}{2} \text{Re} \left\{ \sum_i p_i \vec{\nabla} E_i^* \right\} =$$

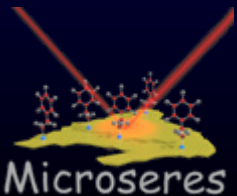
Polarization

$$\langle F \rangle = \frac{1}{4} \Re \{ \alpha \} \nabla |E|^2 + \sigma \frac{1}{2} \Re \left\{ \frac{1}{c} \mathbf{E} \times \mathbf{H}^* \right\} + \sigma \frac{1}{2} \Re \left\{ i \frac{\epsilon_0}{k_0} (\mathbf{E}^* \cdot \nabla) \mathbf{E} \right\}$$

???

Radiation Pressure

$$\frac{1}{c} \langle \vec{S} \rangle \sigma$$



# Optical Forces

(Maxwell eqs. + Lorentz Force)

$$\langle \vec{F} \rangle_t = \frac{1}{2} \text{Re} \left\{ \sum_i p_i \vec{\nabla} E_i^* \right\} =$$

Polarization

$$\langle F \rangle = \frac{1}{4} \Re \{ \alpha \} \nabla |\mathbf{E}|^2 + \sigma \frac{1}{2} \Re \left\{ \frac{1}{c} \mathbf{E} \times \mathbf{H}^* \right\} + \sigma \frac{1}{2} \Re \left\{ i \frac{\epsilon_0}{k_0} (\mathbf{E}^* \cdot \nabla) \mathbf{E} \right\}$$

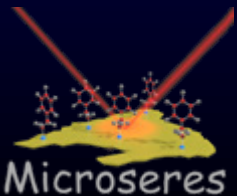
Radiation Pressure

$$\sigma \frac{1}{2} \Re \left\{ i \frac{\epsilon_0}{k_0} (\mathbf{E}^* \cdot \nabla) \mathbf{E} \right\} = \sigma c \nabla \times \left( \frac{\epsilon_0}{4\omega i} (\mathbf{E} \times \mathbf{E}^*) \right)$$

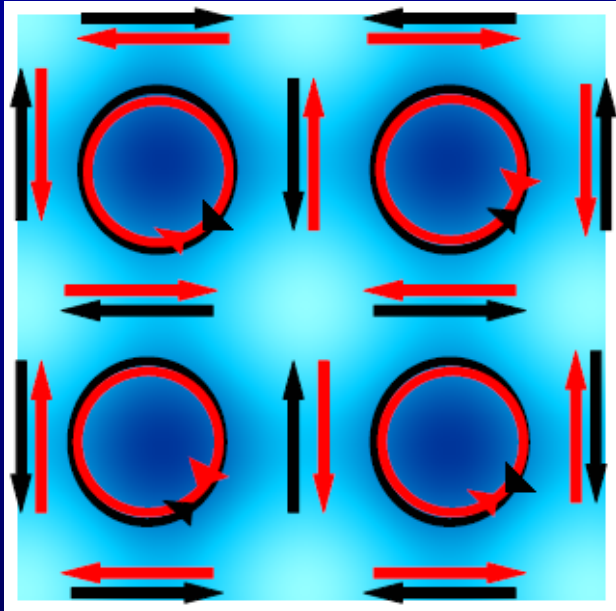
$$\frac{1}{c} \langle \vec{S} \rangle \sigma$$

$$\langle L_S \rangle = \frac{\epsilon_0}{4\omega i} \{ \mathbf{E} \times \mathbf{E}^* \}$$

Time averaged spin density!!



## Optical lattice



$$F^{(p)} = \Re\{\alpha\} \frac{1}{2} |E_0|^2 \nabla (\cos^2 k_0 x + \cos^2 k_0 y).$$

## Poynting vector?

$$\frac{1}{c} \langle S \rangle^{(s)} = c \nabla \times \langle L_O \rangle$$

$$\langle L_O \rangle \equiv \frac{\epsilon_0}{\omega} |E_0|^2 \sin \phi \cos k_0 x \cos k_0 y \mathbf{u}_z$$

$$\langle L_S \rangle = - \langle L_O \rangle \equiv \frac{\epsilon_0}{\omega} |E_0|^2 \sin \phi \cos k_0 x \cos k_0 y \mathbf{u}_z$$

Gradient force

Scattering force

$$\langle \mathbf{F} \rangle = \Re\{\alpha\} \left\{ \nabla \frac{1}{2} \langle |\mathbf{E}|^2 \rangle \right\} + \sigma \left\{ \frac{1}{c} \langle \mathbf{S} \rangle \right\} + \sigma \left\{ c \nabla \times \langle \mathbf{L}_S \rangle \right\}.$$

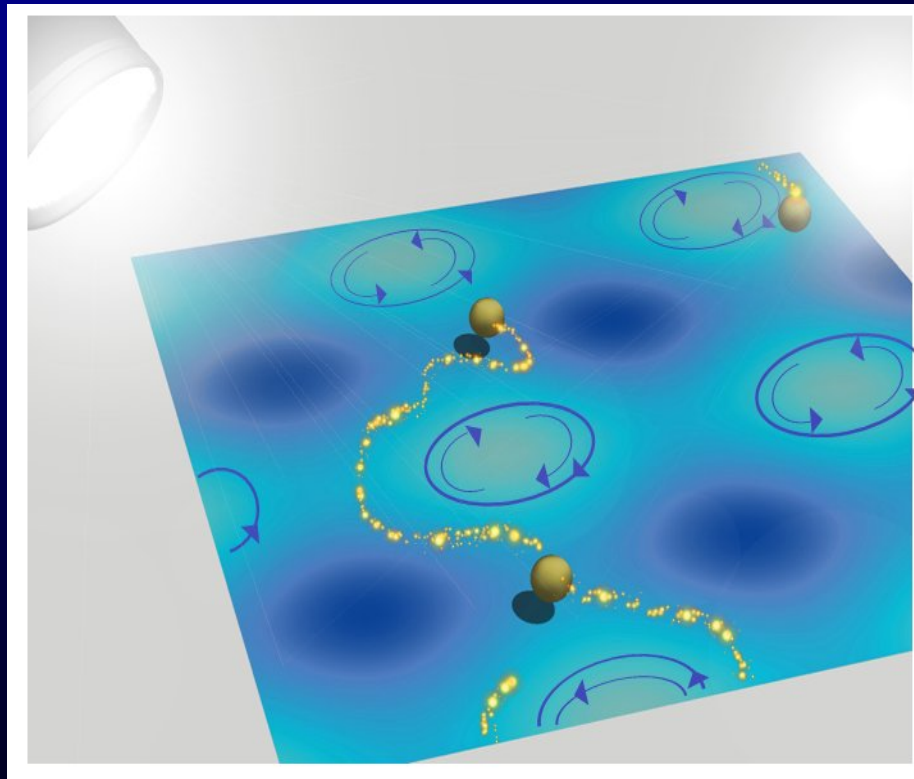
The scattering force (proportional to the total cross section) can be written as the sum of two contributions:

- the traditional radiation pressure term, proportional to the Poynting vector,
- a curl force associated to the non-uniform distribution of the spin density of the light field.

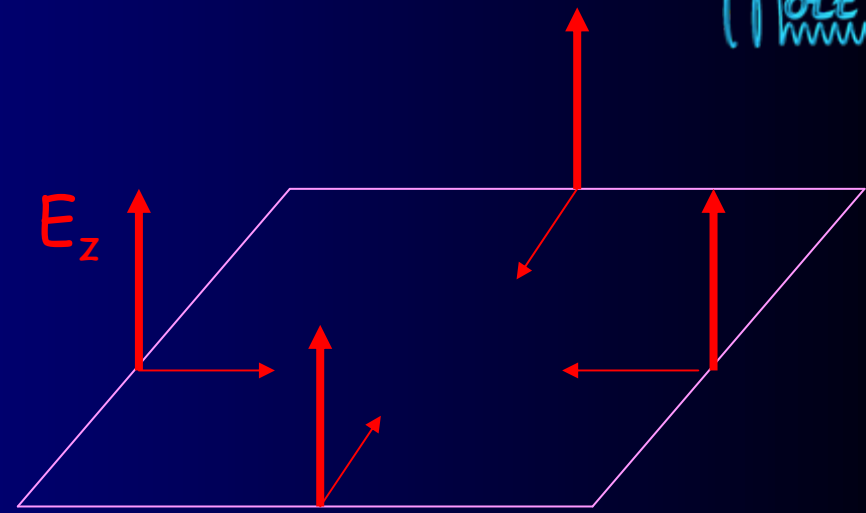
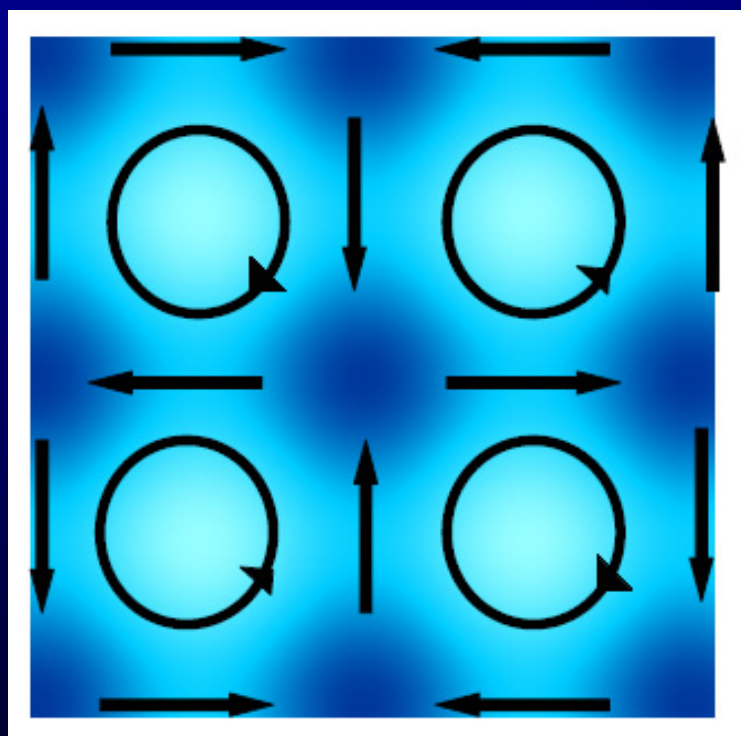
When the light is linearly polarized the curl term is identically zero.

(S. Albaladejo et al., PRL (2009) )

# Giant enhanced diffusion of gold nanoparticles on vortex lattices (S. Albaladejo et al., Nano Letters (2009) )



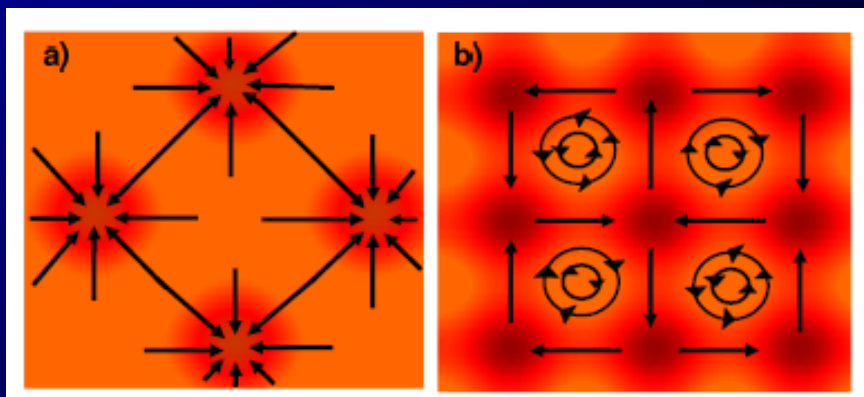
# Optical vortex lattice



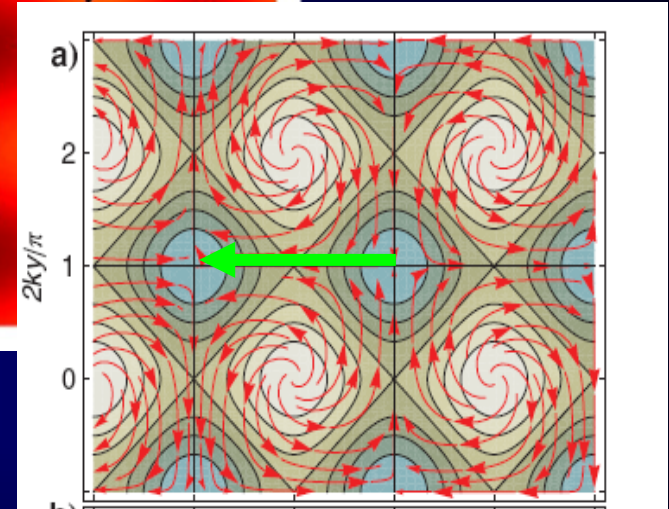
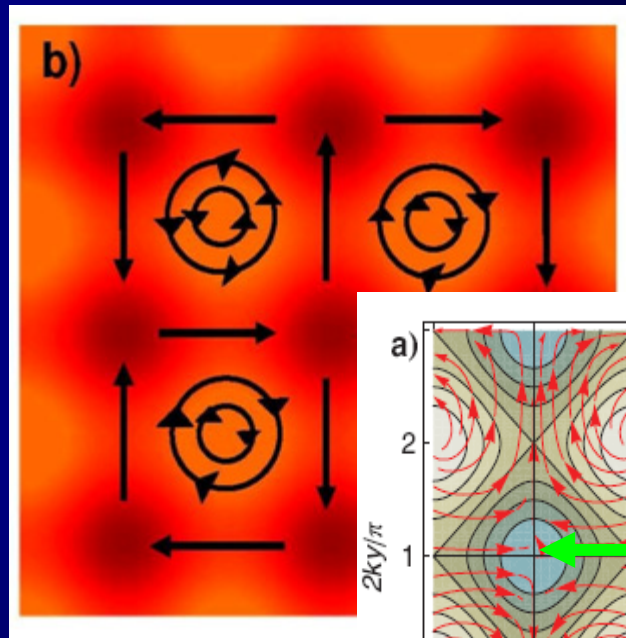
A. Hemmerich and T.W. Hänsch,  
 Phys. Rev. Lett. 68, 1492 (1992)

$$F_{\phi=0} = 2\alpha'(n/c)P\nabla(\sin kx + \sin ky)^2$$

$$(n/c)P \equiv \varepsilon_0\varepsilon|E_0|^2/2.$$



$$F_{\phi=\pi/2} = 2\alpha'(n/c)P\nabla(\sin^2 kx + \sin^2 ky) + 2\alpha''(n/c)P\nabla \times \{2 \cos kx \cos ky u_z\}$$



$$4(n/c)P(\alpha' \sin kx - \alpha'')\cos kx.$$

$$\alpha'' > \alpha',$$

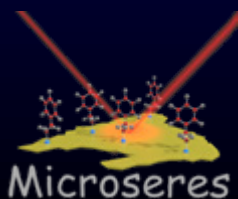
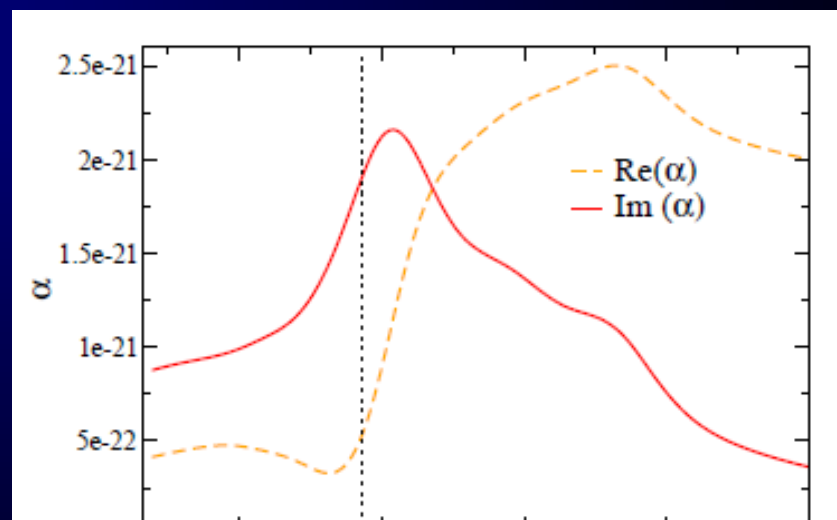
NO Stable positions in  
the system !!!



# Gold nanoparticles

$$\alpha_0 = 4\pi\epsilon_0 a^3 \left( \frac{\epsilon(\omega) - 1}{\epsilon(\omega) + 2} \right)$$

$$\alpha = \frac{\alpha_0}{1 - i \frac{k^3}{6\pi} \alpha_0}$$



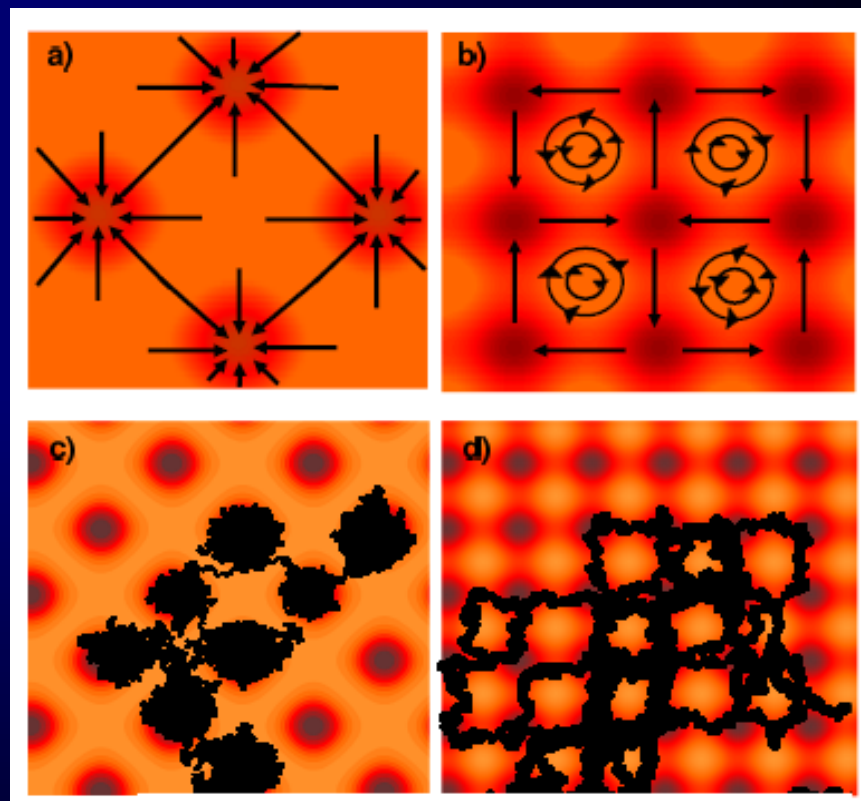
$$m \frac{d^2 r}{dt^2} = F(r) - \gamma \frac{dr}{dt} + \xi(t)$$

$$\gamma = 6\pi a \eta$$

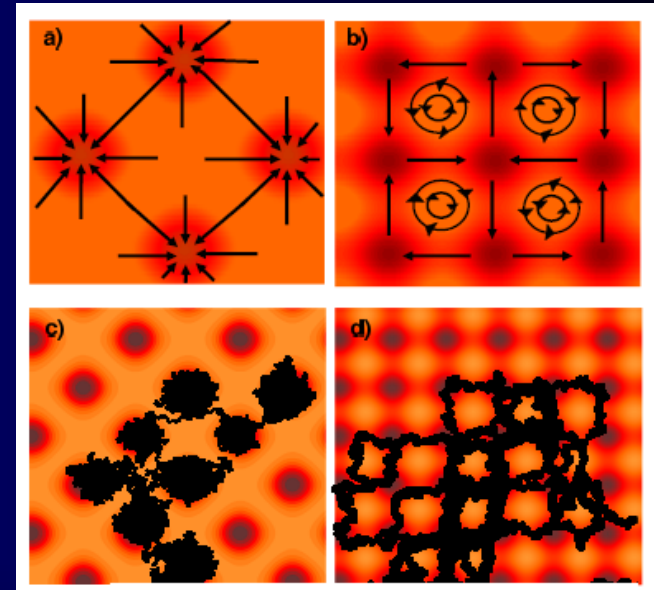
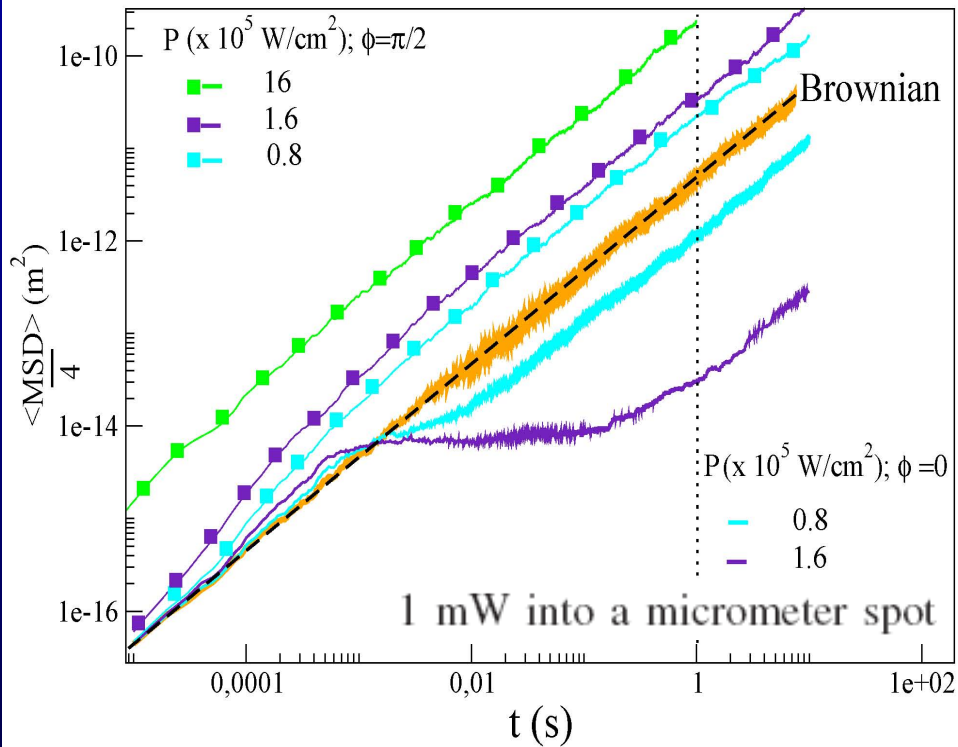
$$T = 298 \text{ K}, \eta = 0.89 \times 10^{-3} \text{ kg m}^{-1} \text{ s}^{-1}$$

$$\langle \xi_i(t) \xi_j(t') \rangle = 2\gamma k_B T \delta_{ij} \delta(t - t')$$

Langevin dynamics



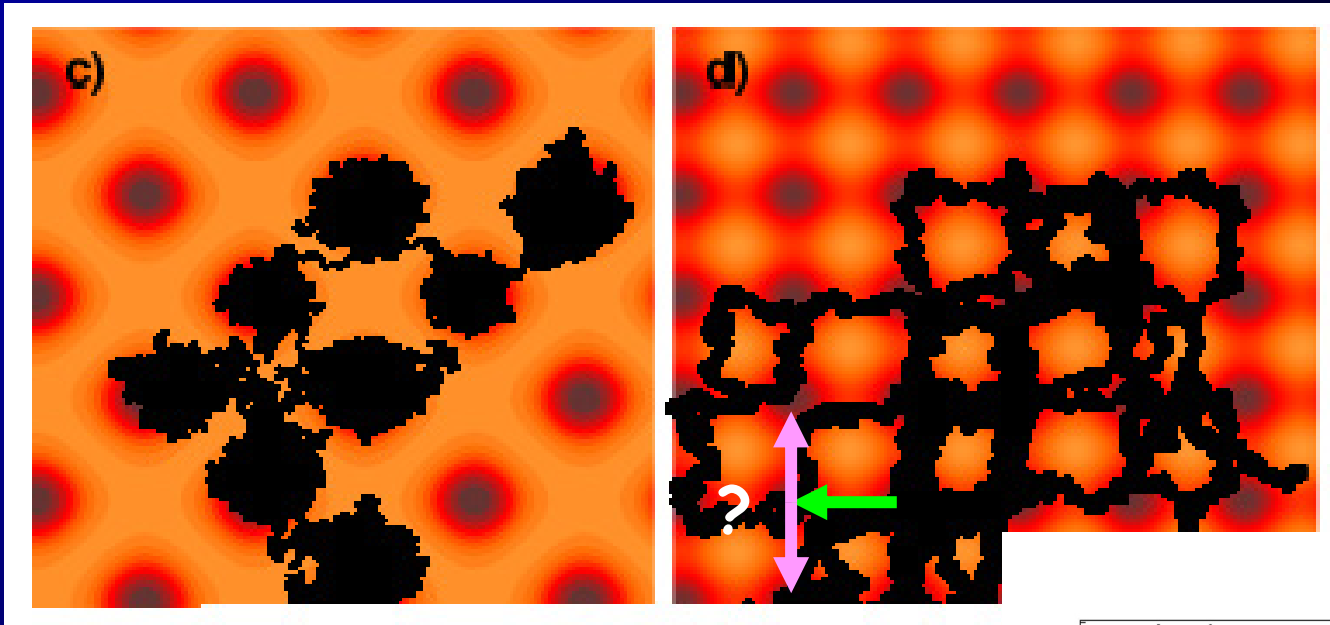
# 50 nm gold nanoparticles in water



$$\langle |x(t) - x(0)|^2 + |y(t) - y(0)|^2 \rangle = \langle r^2 \rangle = 4Dt$$

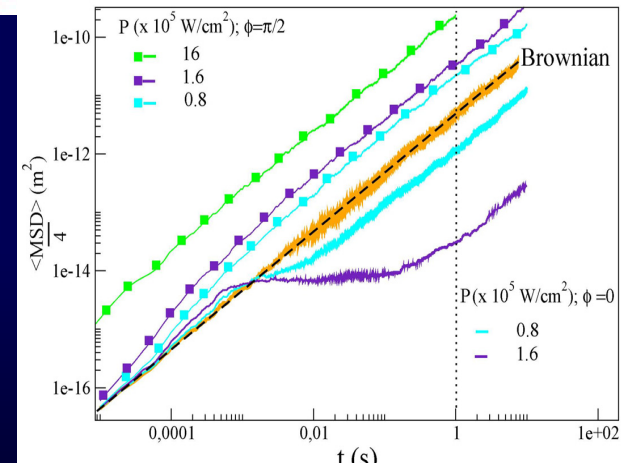
(S. Albaladejo et al.  
NanoLett. 2009)

$$D_0 = k_B T / \gamma \approx 4.9 \times 10^{-12} \text{ m}^2/\text{s}.$$

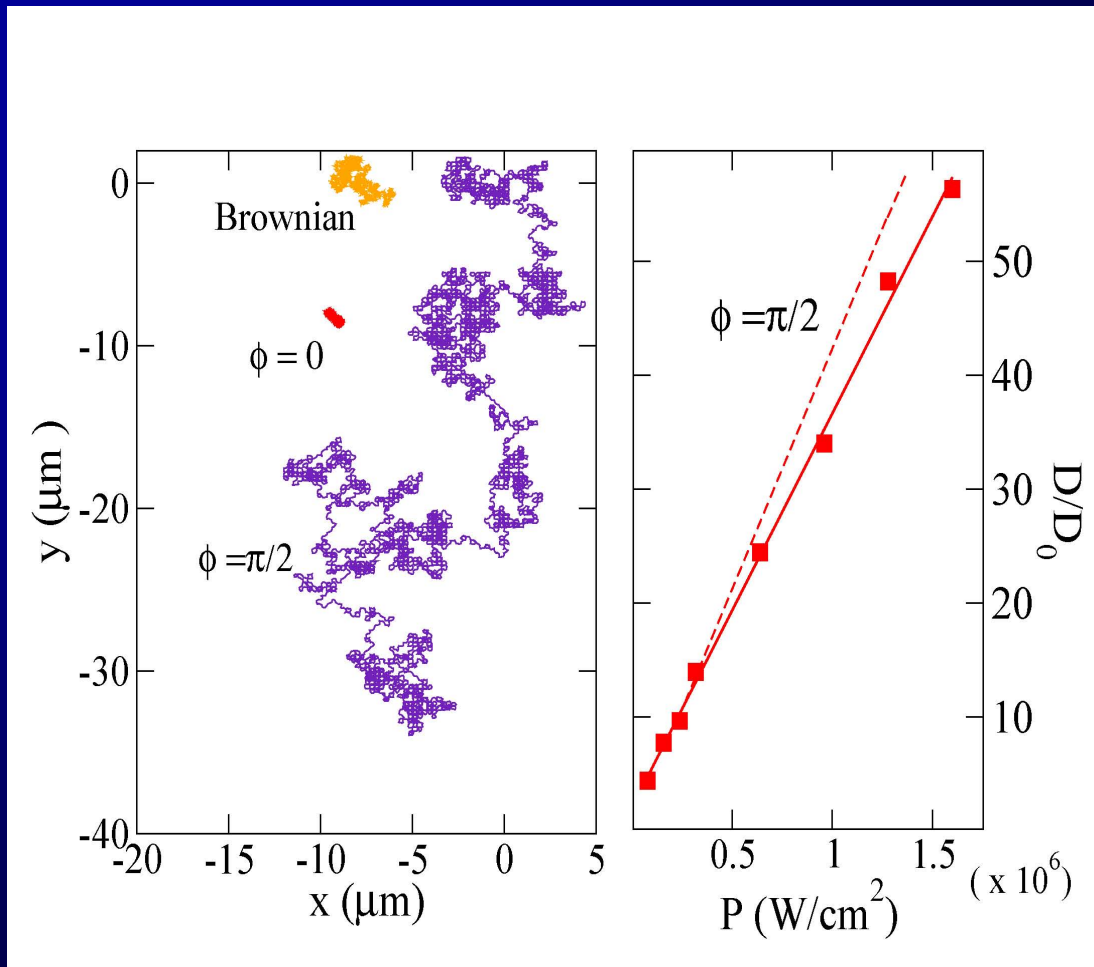


$$\tau \approx 10^{-5} - 10^{-4} \text{ s}$$

$$10^6 - 10^5 \text{ W/cm}^2.$$

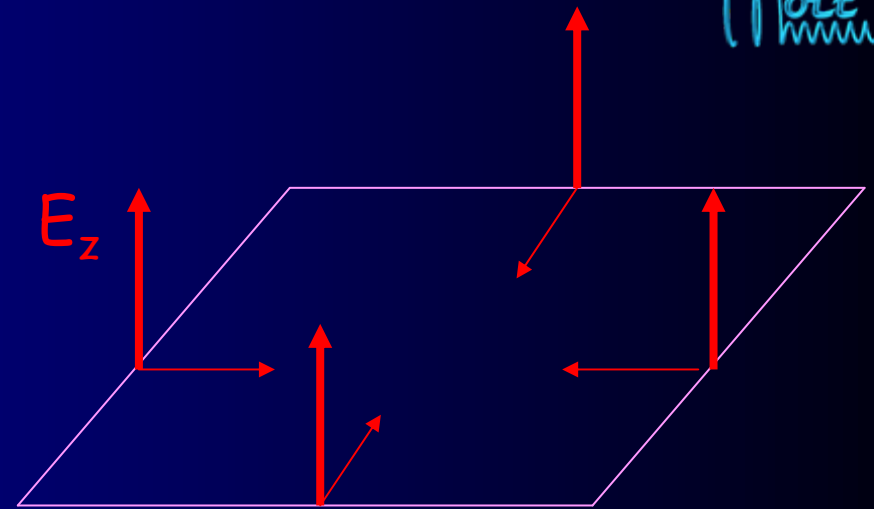
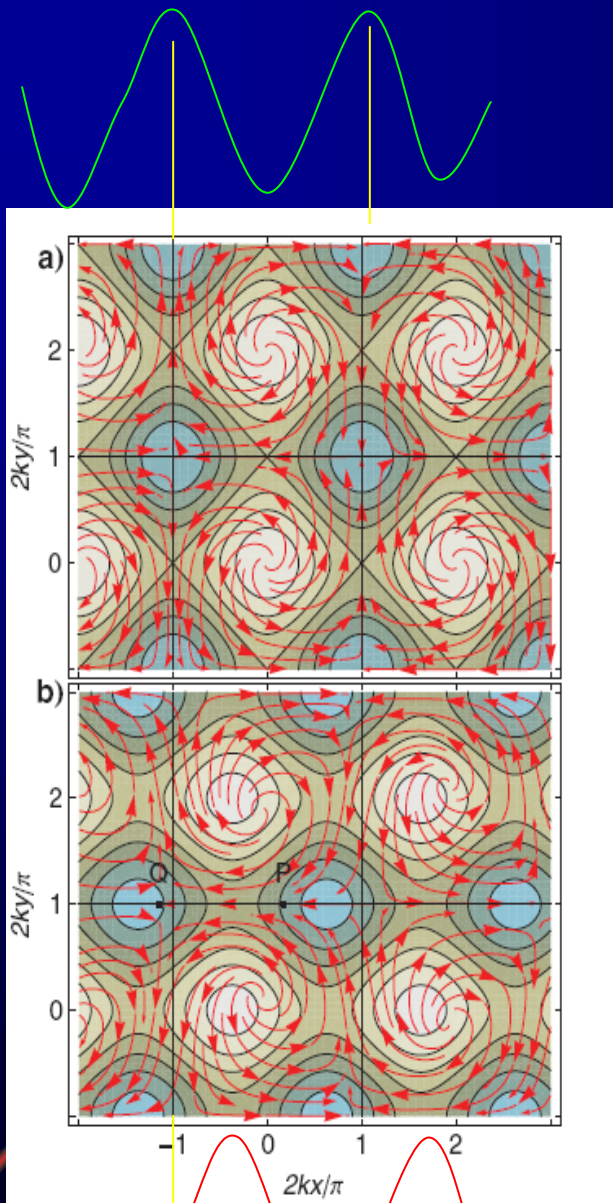


$$D - D_0 \approx \left(\frac{\lambda}{2}\right)^2 \frac{1}{4\tau} \approx 2 \frac{(n/c)\alpha'' P}{\gamma} = D_0 2 \frac{(n/c)\alpha''}{k_B T} P$$

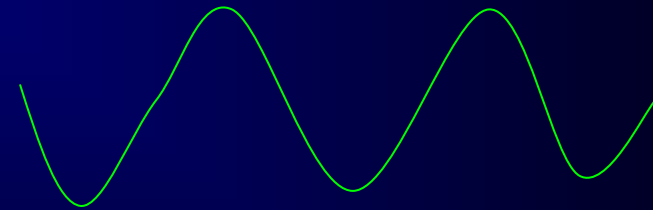


The diffusion constant of a 50 nm gold particle is enhanced by 2 orders of magnitude with respect to thermal diffusion at room temperature

(S. Albaladejo et al. NanoLett. 2009)



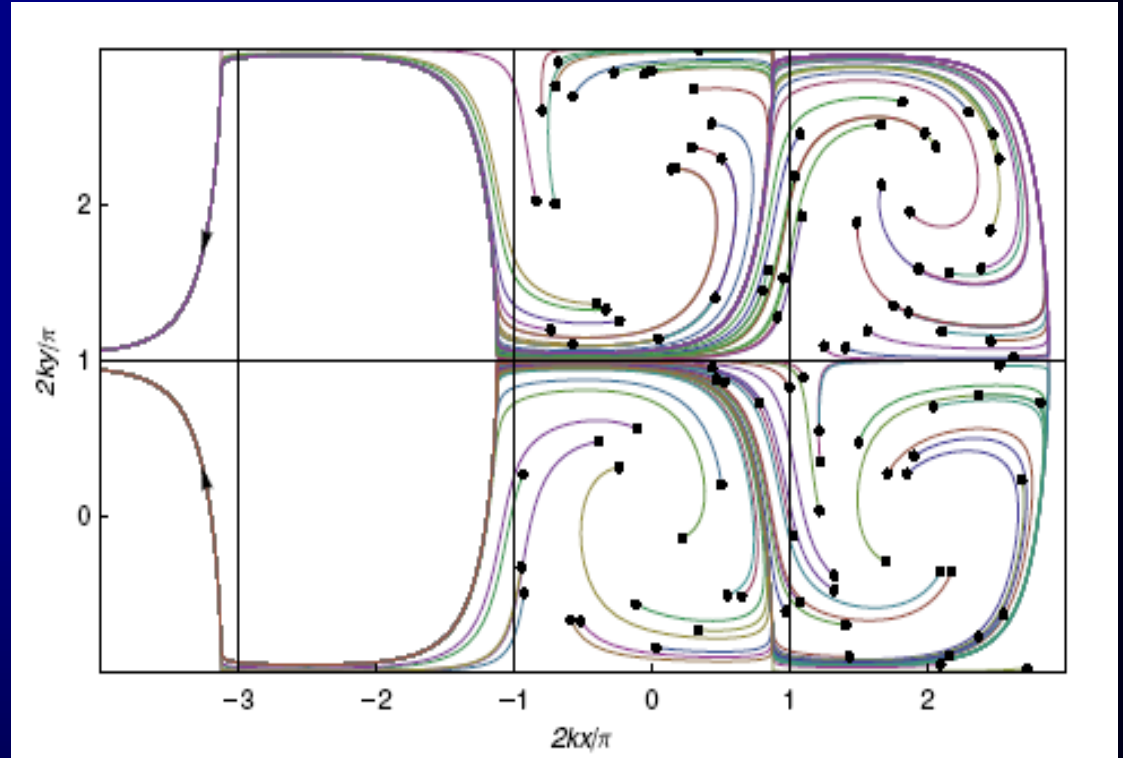
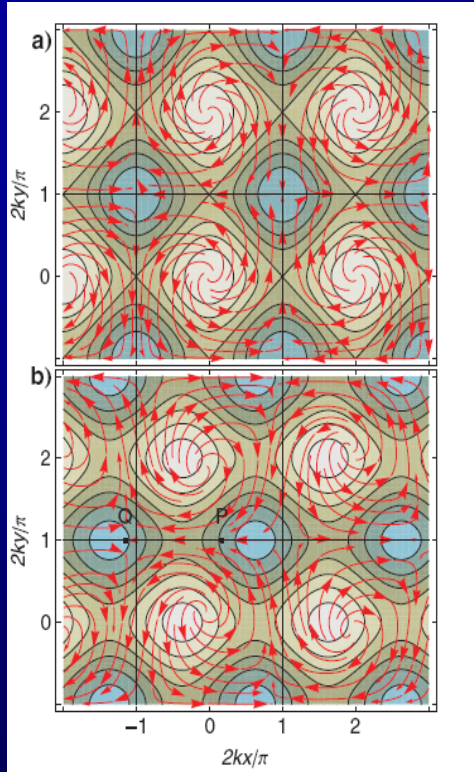
+



## Deterministic Ratchet from Stationary Light Fields

(I. Zapata et al. PRL 2009)





From stationary forces of null average we have designed a ratchet system which requires neither noise nor driving.

(I. Zapata et al. PRL 2009)

