Lateral Ir-spaced superlattices of Fe, Co, and Ni on Ir(100)

A. Schmidt, A. Klein, C. Giovanardi, L. Hammer and K. Heinz

Lehrstuhl für Festkörperphysik, Universität Erlangen-Nürnberg, Staudtstr. 7, D-91058 Erlangen Email: andreas.schmidt@physik.uni-erlangen.de

The stable phase of clean Ir(100) exhibits a quasi-hexagonal top layer with six atomic rows fitting to five rows of the second layer ($Ir(100)-(5\times 1)$ -hex). This leads to a surface buckling which amounts to more than a half of an Angstrøm for the top layer and extends up to the fourth layer into the surface [1]. Exposure to hydrogen lifts this reconstruction by expelling the 20% extra atoms accommodated formerly in the top layer [2]. This $Ir(100)-(5\times 1)$ -H phase is characterized by monoatomic, almost defect free Ir wires of up to micrometer length which reside in 5×1 long-range order on the unreconstructed substrate. As shown in fig. 1 they are almost equally spaced with a lateral distance of 1.36 nm, which corresponds to 5a with a the in-plane lattice parameter (a = 2.72 Å). Yet, also 3a and 7a spacings occur with statistical weight depending severely on preparation conditions. The buckling of the layers beneath the wires is dramatically reduced compared to the 5×1 -hex phase. We show that this surface can be used as a template for lateral superlattices and layer-by-layer growth of magnetic materials. By deposition of Fe, Co or Ni the space between these wires can be filled, whereby the three elements differ by their in-plane growth modes as observed by scanning tunnelling microscopy (STM). Their final states at 0.8 ML coverage correspond to lateral superlattices, {Fe₄Ir}, {Co₄Ir} and {Ni₄Ir} as shown for Fe in fig. 2. The crystallographic structures of these three compounds as determined from the measurement and quantitative dynamical analysis of low-energy electron diffraction (LEED) intensities [3] are very similar, but differ in some parameters. In the case of iron, the adatoms being next to the Ir rows are shifted by about 0.05 Å towards Ir, while Co and Ni atoms experience an only small shift in the other direction. This indicates a larger preference for homogeneous bonds for Ni and Co in contrast to Fe. For all three types adatoms intermixing with Iridium can be excluded.

The Fe₄Ir surface superlattice can also be obtained without the use of hydrogen as surfactant [4]. The deposition of 0.8 ML of Fe on Ir(100)-(5×1)-hex makes the surface restructure in the same way and a Fe₄Ir compound with the same local geometry is formed. Yet, this preparation procedure leads to a poor degree of long-range order.

Also, Fe-, Co- and Ni-films of higher coverages on the $Ir(100)-(5\times1)$ -H template are analysed. The growth is found to be layer-by-layer. For e.g. a deposition of 1.8 ML Ni, the Ni₄Ir superlattice is fully covered by pure Ni, with the chain like structure of the interface, preserved, as demonstrated in fig. 3. Figure 1: STM image of the restructured Ir(100)-(5×1) after exposure to hydrogen at room temperature. The small picture shows the monoatomic Ir wires with a distance of 5*a*.

Figure 2: Left: STM image after deposition of approx 0.8 ML Fe on the Ir template (note that the Ir wires appear as dark lines). Right: Ball model of the Fe₄Ir superlattice on a STM image of the monoatomic Ir wires for visualization.

Figure 3: STM image after the deposition of 1.8 ML Ni on the Ir(100)-H showing the preserved chain like structure of the interface.

References

| [1] | A. Schmidt, W. Meier, L. Hammer and K. Heinz, J. Phys Condens. Matter 14 (2002) 12353 |
|-----|---------------------------------------------------------------------------------------------------------------|
| [2] | L. Hammer, W. Meier, A. Klein, P. Landfried, A. Schmidt and K. Heinz, Phys. Rev. Lett. 14 (2002) 12353 |
| [3] | V. Blum and K. Heinz, Comp. Phys. Comm. 134 2001 392 |
| [4] | L. Hammer, W. Meier, A. Schmidt and K. Heinz, Phys. Rev. B. 67 (2003) 125422 |



